

1.1. Differential forms.

a) Show that a closed n -form ω on the unit sphere $S^n \subset \mathbb{R}^{n+1}$ is exact if and only if $\int_{S^n} \omega = 0$.

b) Prove that a compactly supported n -form ω on \mathbb{R}^n is the exterior derivative of a compactly supported $(n-1)$ -form on \mathbb{R}^n if and only if $\int_{\mathbb{R}^n} \omega = 0$.

Hint: Use stereographic projection to carry ω to a form ω' on S^n . By part a), $\omega' = d\nu$. Now $d\nu$ is 0 in a contractible neighborhood U of the north pole N . Use this to find an $n-2$ form μ on S^n such that $\nu = d\mu$ near N . Then $\nu - d\mu$ is zero near N , so it pulls back to a compactly supported form on \mathbb{R}^n .

Solution. a) By Stokes' Theorem, integration yields a well-defined linear map

$$I: H_{\text{dR}}^n(S^n) \rightarrow \mathbb{R}.$$

Since the integral of a volume form is non-zero, I is surjective. Furthermore, as seen in class, $H_{\text{dR}}^n(S^n)$ is 1-dimensional, so I is an isomorphism. The kernel, the class $[0]$, is represented precisely by the exact forms.

b) The hint gives essentially the full solution. □

1.2. De Rham cohomology of T^2 . Determine the De Rham cohomology of the torus T^2 .

Solution. Cover the torus T^2 using two (sets homeomorphic to) cylinders U, V so that they overlap slightly at the extremities. Then U and V are both homotopy equivalent to S^1 , so

$$H_{\text{dR}}^s(U) \cong H_{\text{dR}}^s(V) \cong H_{\text{dR}}^s(S^1) \cong \begin{cases} \mathbb{R} & s = 0, 1, \\ 0 & s \geq 2. \end{cases}$$

The intersection $U \cap V$ is homotopy equivalent to the disjoint union of two copies of S^1 and it is not difficult to see (e.g. by using the Mayer-Vietoris sequence) that

$$H_{\text{dR}}^s(U \cap V) \cong H_{\text{dR}}^s(S^1 \sqcup S^1) \cong H_{\text{dR}}^s(S^1) \oplus H_{\text{dR}}^s(S^1) \cong \begin{cases} \mathbb{R}^2 & s = 0, 1, \\ 0 & s \geq 2. \end{cases}$$

Using that $H_{\text{dR}}^0(T^2) \cong \mathbb{R}$ and, as seen in the lecture, $H_{\text{dR}}^2(T^2) \cong \mathbb{R}$, the Mayer-Vietoris sequence for this open cover is

$$0 \rightarrow H_{\text{dR}}^0(T^2) \rightarrow H_{\text{dR}}^0(U) \oplus H_{\text{dR}}^0(V) \rightarrow H_{\text{dR}}^0(U \cap V) \rightarrow H_{\text{dR}}^1(T^2)$$

$$\rightarrow H_{dR}^1(U) \oplus H_{dR}^1(V) \rightarrow H_{dR}^1(U \cap V) \rightarrow H_{dR}^2(T^2) \rightarrow 0,$$

which amounts to

$$0 \rightarrow \mathbb{R} \rightarrow \mathbb{R}^2 \rightarrow \mathbb{R}^2 \rightarrow H_{dR}^1(T^2) \rightarrow \mathbb{R}^2 \rightarrow \mathbb{R}^2 \rightarrow \mathbb{R} \rightarrow 0.$$

Since the alternating sum of the dimensions is 0, we conclude that $-1 + 2 - 2 + \dim H_{dR}^1(T^2) - 1 + 2 - 1 = 0$, and therefore $H_{dR}^1(T^2) \cong \mathbb{R}^2$. \square

1.3. Tensor fields. Let T be a $(1,2)$ -tensor field on M^m . Let (φ, U) and (ψ, U) be two charts on M . Show that the component ${}^\psi T_{ab}^c$ of T with respect to ψ depends on the components ${}^\varphi T_{ij}^k$ of T with respect to φ by the following relation:

$${}^\psi T_{ab}^c = \sum_{i,j,k=1}^m \frac{\partial \psi^c}{\partial \varphi^k} \frac{\partial \varphi^i}{\partial \psi^a} \frac{\partial \varphi^j}{\partial \psi^b} {}^\varphi T_{ij}^k.$$

Solution. For $a, b, c \in \{1, \dots, m\}$, we have

$${}^\psi T_{ab}^c = T(d\psi^c \otimes \frac{\partial}{\partial \psi^a} \otimes \frac{\partial}{\partial \psi^b})$$

so by writing T with respect to φ we see that

$$\begin{aligned} {}^\psi T_{ab}^c &= \sum_{i,j,k=1}^m {}^\varphi T_{ij}^k \frac{\partial}{\partial \varphi^k} \otimes d\varphi^i \otimes d\varphi^j \left(d\psi^c \otimes \frac{\partial}{\partial \psi^a} \otimes \frac{\partial}{\partial \psi^b} \right) \\ &= \sum_{i,j,k=1}^m {}^\varphi T_{ij}^k \left(d\psi^c \left(\frac{\partial}{\partial \varphi^k} \right) \right) \left(d\varphi^i \left(\frac{\partial}{\partial \psi^a} \right) \right) \left(d\varphi^j \left(\frac{\partial}{\partial \psi^b} \right) \right). \end{aligned}$$

Since

$$d\psi^c \left(\frac{\partial}{\partial \varphi^k} \right) = \frac{\partial \psi^c}{\partial \varphi^k}, \quad d\varphi^i \left(\frac{\partial}{\partial \psi^a} \right) = \frac{\partial \varphi^i}{\partial \psi^a}, \quad d\varphi^j \left(\frac{\partial}{\partial \psi^b} \right) = \frac{\partial \varphi^j}{\partial \psi^b},$$

substituting these into the sum gives

$${}^\psi T_{ab}^c = \sum_{i,j,k=1}^m {}^\varphi T_{ij}^k \frac{\partial \psi^c}{\partial \varphi^k} \frac{\partial \varphi^i}{\partial \psi^a} \frac{\partial \varphi^j}{\partial \psi^b}.$$

\square